

# Experimental Status of B-physics

Hitoshi Yamamoto

University of Hawaii

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## 1. B-factories

## 2. QCD and Weak Interaction

- Inclusive hadronic rates
- Radiative  $B$  decays ( $b \rightarrow s\gamma$ )

## 3. CP violation and Unitarity Triangle

- $|V_{ub}|$  and  $|V_{cb}|$  ( $|V_{td}|$ )
- CP phase  $\beta$  ( $\phi_1$ )
- Direct  $CP$  violations

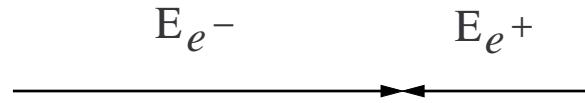
# B-factories

$< e^+e^- \text{ B-factory accelerators} >$

## Symmetric energies (CESR)

$$E_{e^-} = E_{e^+} = \frac{M_{\gamma 4S}}{2} = 5.29 GeV$$

## Asymmetric energies (PEP-II, KEK-B)



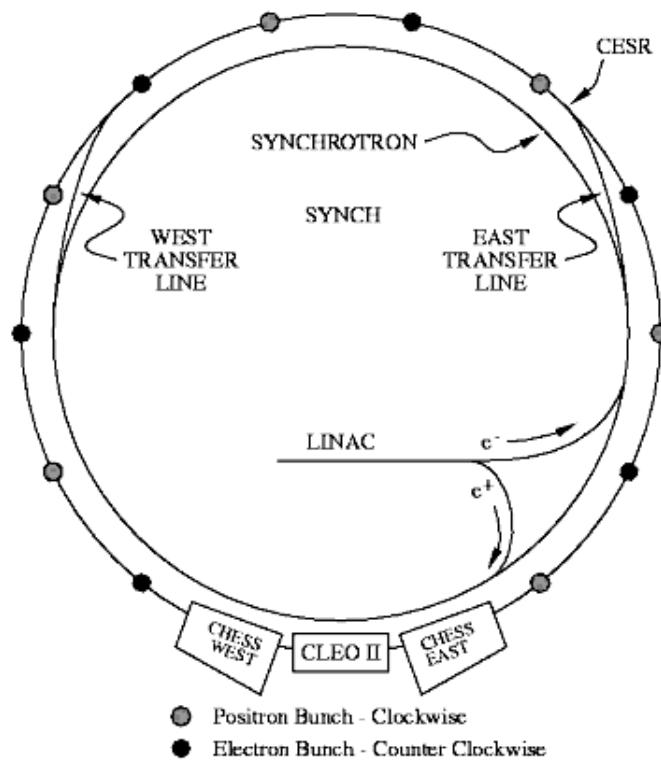
$\gamma 4S$  is moving in the lab frame.

$$E_{CM} = 2\sqrt{E_{e^+}E_{e^-}} = M_{\gamma 4S}$$

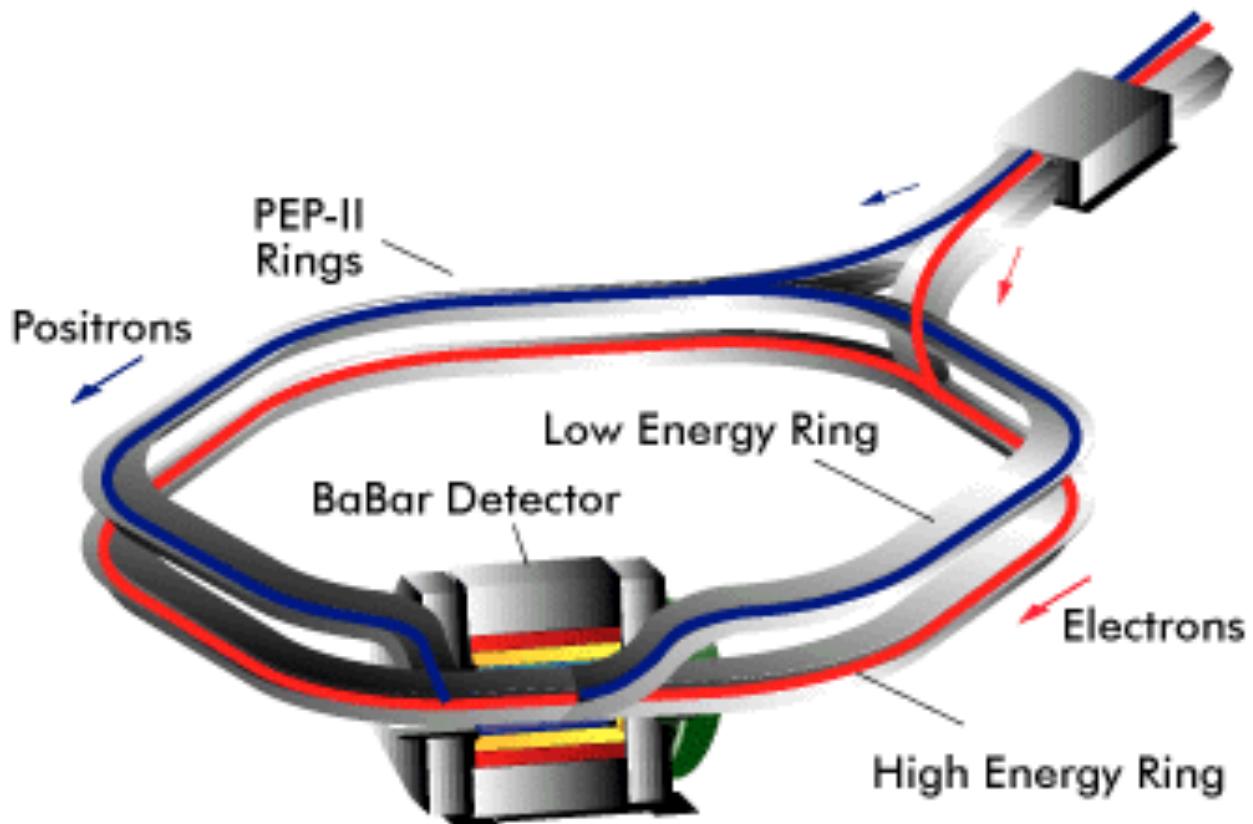
$$\begin{cases} E_{\gamma 4S} = E_{e^-} + E_{e^+} \\ P_{\gamma 4S} = E_{e^-} - E_{e^+} \end{cases}$$

$$\rightarrow \beta_{\gamma 4S} = \frac{P_{\gamma 4S}}{E_{\gamma 4S}} = \frac{E_{e^-} - E_{e^+}}{E_{e^-} + E_{e^+}}$$

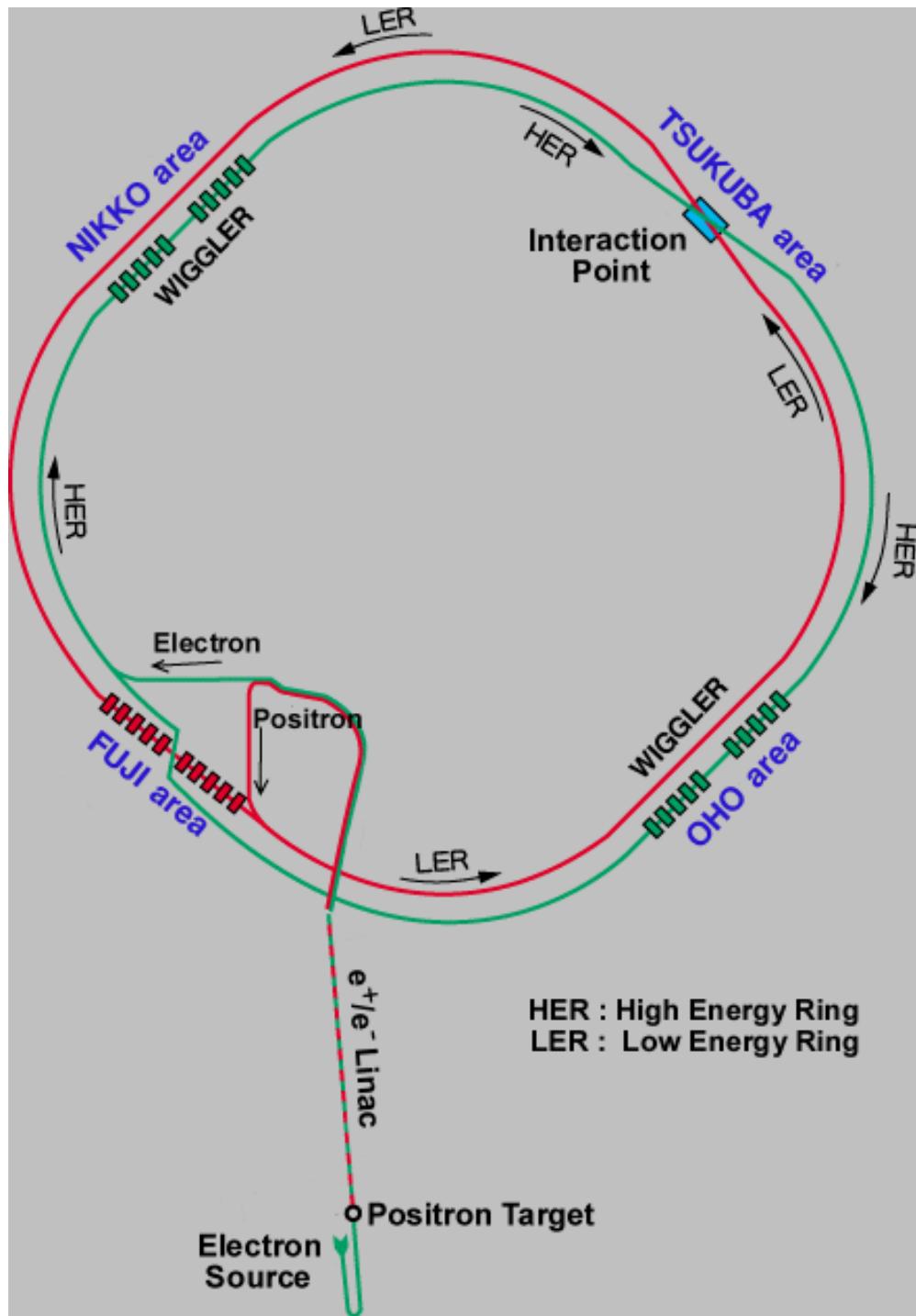
## CESR (Cornell Electron Strange Ring)



## PEP-II (SLAC)



## KEK-B (KEK, Japan)



## Beam separation

Want collision to occur only at one location

→ Need for beam separation  
(avoid parasitic crossings)

**CESR:** Pretzel orbit

Interweaving  $e^+e^-$  orbits within a single ring  
Crossing angle =  $\pm 2.3$  mrad

**PEP-II:** Separation by bending magnet

$$E_{e^+} \neq E_{e^-}$$

→  $e^+, e^-$  beams bend differently

Head-on collision

**KEK-B:** Finite-angle crossing

Crossing angle =  $\pm 11$  mrad

Large crossing angle

Beam instability

(synchro-betatron resonance)

→ Luminosity reduction  
(geometrical)

machine	<i>CESR</i>	<i>PEP-II</i>	<i>KEK-B</i>
detector	<i>CLEO</i>	<i>BaBar</i>	<i>Belle</i>
circumference (km)	0.768	2.199	3.016
# of rings	1	2	2
$E_{e^+}$ (GeV)	5.3	3.1	3.5
$E_{e^-}$ (GeV)	5.3	9.0	8.0
$\beta\gamma_{4S}$	~ 0	0.49	0.39
$\delta E/E$	$6 \times 10^{-4}$	$7 \times 10^{-4}$	$7 \times 10^{-4}$
$\Delta t_{\text{bunch}}$	14ns	4.2ns	2ns
bunch size( $w$ )	500 $\mu$	181 $\mu$	77 $\mu$
" (h)	10 $\mu$	5.4 $\mu$	1.9 $\mu$
" (l)	1.8cm	1.0cm	0.4cm
crossing angle(mrad)	$\pm 2.3$	0	$\pm 11$
Luminosity( $cm^{-2}s^{-1}$ )	$1.5 \times 10^{33}$	$3 \times 10^{33}$	$10^{34}$
# $B\bar{B}/s$	1.5	3	10

achievements so far

Lum(peak)	$8 \times 10^{32}$	$1.5 \times 10^{33}$	$6 \times 10^{32}$
$\int Ldt$ (fb $^{-1}$ )	9.2	1.7	0.3

## B-factory Status

- CLEO-3 detector currently running with a temporary vertex detector.  
To be completed in Feb 2000.
- BaBar and Belle currently running.  
 $\sin 2\beta$  result in summer 2000??

## Hadronic machines

- HERA-B: wire target inside  $e^- - p$  collider at DESY.  
HERA-B begin assembled.  
To be completed Feb 2000.
- CDF/D0: Fermilab  $p\bar{p}$  collider.  
Run2 to begin in fall 2000.

## Future machines

- LHCb, BTeV, ATLAS, CMS:  $\sim 2006$  and on.

**LEP** Finished  $Z^0$  runs - no more  $B$ -physics data.  
Some  $B$  results still to come on the existing data.

# Inclusive Hadronic Rates

Two parameters:

1.  $B_{\text{s.l.}} \equiv \text{average } Br(H_b \rightarrow X e^- \nu_e)$

where  $H_b = \begin{cases} \bar{B}^0, B^- & (\Upsilon 4S) \\ \bar{B}^0, B^-, \bar{B}_s^0, N_b(\text{any } b\text{-baryon}) & (Z^0) \end{cases}$

Expect  $B_{\text{s.l.}}(\Upsilon 4S) = \frac{\tau_B}{\tau_b} B_{\text{s.l.}}(Z^0)$

$$B_{\text{s.l.}} = \frac{1}{2.22 + r_{\text{had}}}, \quad r_{\text{had}} = r_{c\bar{u}d} + r_{c\bar{c}s} + r_{\text{noc}}$$

( $r_x$  : rate in unit of  $\Gamma_{\text{s.l.}}$ )

$$2.22 = 1(e) + 1(\mu) + 0.22(\tau)$$

2.  $n_c \equiv \text{number of } c \text{ or } \bar{c} \text{ per } B \text{ decay}$

Weakly decaying charms:  $D^0, D^+, D_s^+, \Lambda_c, \Xi_c$

Charmonium  $\equiv 2c$  if decayed by  $c\bar{c}$  annihilation.

Expect more  $\Lambda_c, \Xi_c$ , and  $D_s$  on  $Z^0$ .

$$n_c = 1 + Br_{c\bar{c}s} - Br_{\text{noc}} = 1 + \frac{r_{c\bar{c}s} - r_{\text{noc}}}{2.22 + r_{\text{had}}}$$

## Theoretical estimates of $B_{\text{s.l.}}$ and $n_c$

### Tools

- Local quark-hadron duality

$\Sigma$  of final states with certain flavor contents  
=  $\Sigma$  of the same flavor contents at quark level

Assumed to hold at fixed  $\sqrt{s}$  of hadronic system  
(‘local’ duality)

- pQCD and Heavy-Quark Expansion (HQE)

$$\Gamma(B \rightarrow f) = \Gamma_0 \left[ \textcolor{violet}{a} \left( 1 + \frac{\lambda_1}{2m_b^2} + \frac{3\lambda_2}{2m_b^2} \right) + \textcolor{violet}{b} \frac{\lambda_2}{m_b^2} + \mathcal{O}\left(\frac{1}{m_b^3}\right) \right]$$
$$\left( \Gamma_0 = \frac{G_F^2 m_b^5}{192\pi^3} \right)$$

$a, b$ : contains CKM factors, Wilson coefficients (pQCD),  
and phase space factors.

$\lambda_{1,2}$ : ‘non-perturbative’ effects

$\lambda_1 = 0 \sim -0.7 \text{ GeV}^2$ : time-delation by fermi motion

$\lambda_2 = \frac{1}{4}(m_B^{*2} - m_B^2) = 0.12 \text{ GeV}^2$   
: chromomagnetic effect

Total effect by  $\lambda_{1,2} \sim$  a few % increase.

Not all non-perturbative effects are included in  $\lambda_{1,2}$ .

Effects of various corrections affecting  $B_{\text{S.I.}}$  and  $n_c$ .  
 (NLO calc by Bagan et. al.)

	naive	NLO( $m_c = 0$ )	NLO( $m_c \neq 0$ )
$r_{cl\nu}$	2.22	2.22	2.22
$r_{c\bar{u}d}$	3.0	4.0	4.0
$r_{c\bar{c}s}$	1.2	1.6	2.1
$B_{\text{S.I.}}$	0.16	0.13	0.12
$n_c$	1.16	1.17	1.21
$Br_{c\bar{c}s}$	0.18	0.20	0.25

$(r_{\text{no c}} \sim 0.2 \pm 0.1)$

Theoretical estimate of  $B_{\text{S.I.}}$  strongly depends on the renormalization scale (also on  $m_c/m_b$ ).

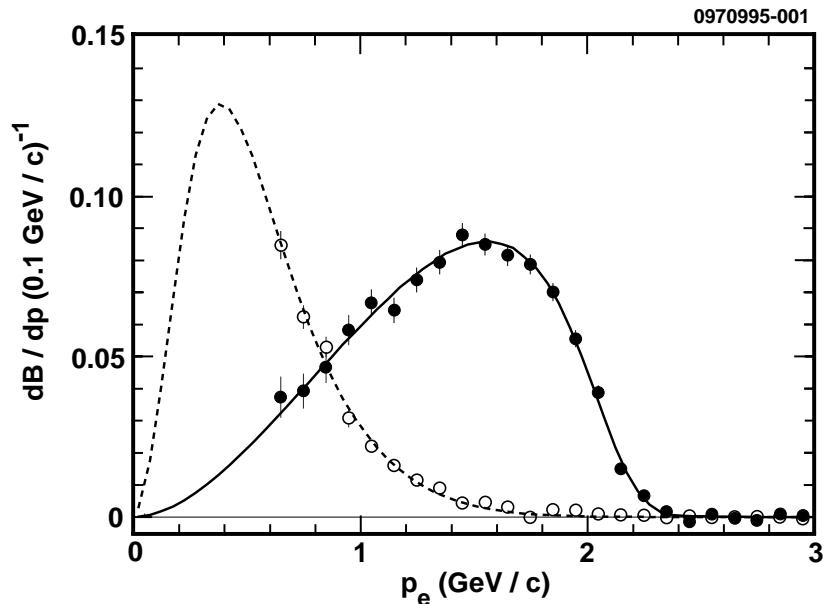
(Neubert, Sachrajda)		
$\mu/m_b$	1	1/2
$B_{\text{S.I.}}(\%)$	$12.0 \pm 1.0$	$10.9 \pm 1.0$
$n_c$	$1.20 \pm 0.06$	$1.21 \pm 0.06$

$$(m_c/m_b = 0.29)$$

# Measurements of $B_{\text{s.l.}}$ .

On  $\Upsilon 4S$  (CLEO) - charge correlation method

- Tag one  $B$  by high- $P$  lepton (flavor-tagged)  
‘The other side’: select  $b \rightarrow e^-$ , reject  $b \rightarrow c \rightarrow e^+$ .



- Correction for  $B^0$ - $\bar{B}^0$  mixing  
Correction for  $b \rightarrow \bar{c} \rightarrow e^-$  ( $B \rightarrow D\bar{D}X$  etc.)

$$B_{\text{s.l.}} = 10.49 \pm 0.46\% \quad (\text{CLEO}).$$

$$B_{\text{s.l.}} = 10.45 \pm 0.21\% \quad (\Upsilon 4S, \text{PDG})$$

## $B_{\text{s.l.}}$ on $Z^0$ (LEP)

Experiment	$B_{\text{s.l.}} (\%)$	method
ALEPH 95	$11.01 \pm 0.10 \pm 0.30$	$1\ell + 2\ell$
L3 96	$10.68 \pm 0.11 \pm 0.46$	$e, \mu, \nu$
OPAL 99	$10.83 \pm 0.10 \pm 0.20^{+0.20}_{-0.13}$	(a)
DELPHI 99	$10.65 \pm 0.07 \pm 0.25^{+0.28}_{-0.12}$	(b)

(a)  $b$ -tag by lifetime.

Fit two neural-net variables to separate  $b \rightarrow \ell^-$ ,  $b \rightarrow c \rightarrow \ell^+$ , and backgrounds.

(b) 4 methods

1.  $1\ell + 2\ell$
2.  $b$ -tag by lifetime/lepton, Fit  $p_\ell(B$  c.m.) and tag-lepton charge correlation.
3. Use all hadronic events + multi-variate fit
4. Similar to 2. Neural net for flavor id.

$$B_{\text{s.l.}} = 10.79 \pm 0.17\% \quad (Z^0 \text{ average}).$$

## Measurement of $n_c$

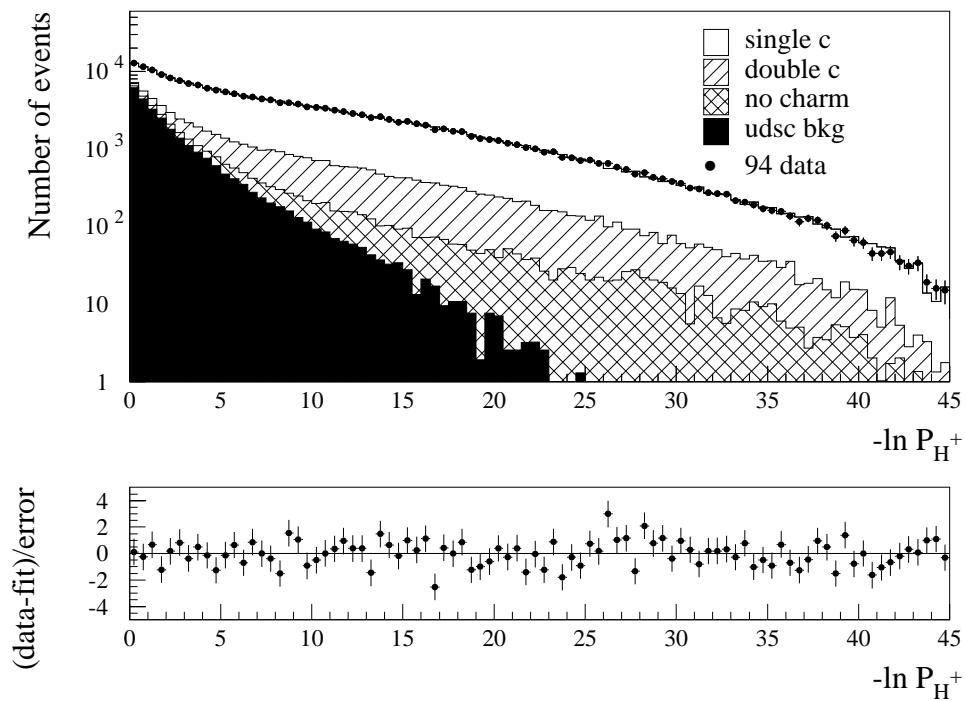
(%)	CLEO 96	ALEPH 96	OPAL 96	DELPHI 99
$D^0$	$63.6 \pm 3.0$	$60.5 \pm 3.6$	$53.5 \pm 4.1$	$60.05 \pm 4.29$
$D^+$	$23.5 \pm 2.7$	$23.4 \pm 1.6$	$18.8 \pm 2.0$	$23.01 \pm 2.13$
$D_s^+$	$11.8 \pm 1.7$	$18.3 \pm 5.0$	$20.8 \pm 3.0$	$16.65 \pm 4.50$
$\Lambda_c$	$3.9 \pm 2.0$	$11.0 \pm 2.1$	$12.5 \pm 2.6$	$8.90 \pm 3.00$
$\Xi_c^{0,+}$	$2.0 \pm 1.0$	$6.3 \pm 2.1$	—	$4.00 \pm 1.60$
$(c\bar{c})$	$5.4 \pm 0.7$	$3.4 \pm 2.4$	—	$4.00 \pm 1.29$
$n_c$	$1.10 \pm 0.05$	$1.23 \pm 0.07$		$1.17 \pm 0.09$

- $\Xi_c$ : ALEPH and DELPHI used the CLEO measurement of  $Br(B \rightarrow \Xi_c)$  and added  $Br(\Lambda_b \rightarrow \Xi_c)$  prediction by JETSET.  
ALEPH used an old (wrong) CLEO number for  $Br(B \rightarrow \Xi_c)$  ( $n_c : 1.23 \rightarrow 1.21$ )
- DELPHI values for  $\Xi_c^{0,+}$  and  $(c\bar{c})$  + OPAL  
 $\rightarrow n_c = 1.14 \pm 0.08$

## Charm counting by vertexing (DELPHI)

Does not depend on charm decay Br's.

Fit  $P_H^+$  (probability that all tracks with positive lifetime comes from the primary vertex)



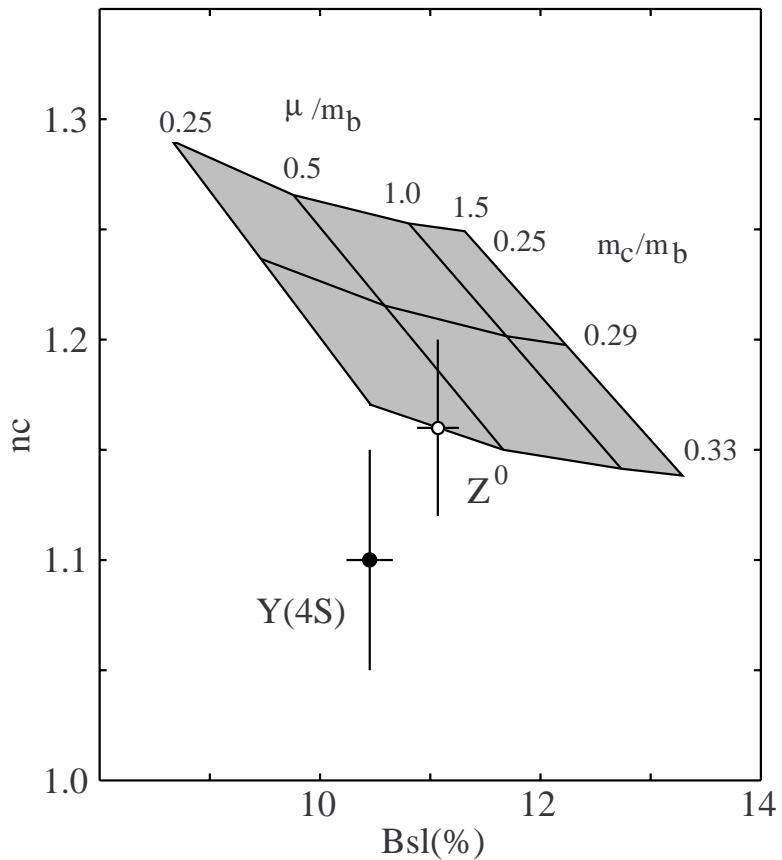
Find numbers of 0,1,2-charm decays

$$n_c = 1.147 \pm 0.041 \pm 0.008 \quad (DELPHI)$$

LEP average:

$$n_c = 1.16 \pm 0.04 \quad (Z^0)$$

## $B_{\text{S.I.}}$ vs $n_c$ . Theory vs experiments.



1.  $B_{\text{S.I.}}$  for  $Z^0$  is converted to  $\Upsilon 4S$  value ( $\times \tau_B / \tau_b$ )
2. Discrepancy between theory and  $\Upsilon 4S$  values is alarming. Boosting  $r_{n_c}$  will fix it.
3. Theory and  $Z^0$  values are consistent.

# Radiative $B$ Decays ( $b \rightarrow s\gamma$ )

## Triumph of OPE/pQCD

Relevant Hamiltonian for  $b \rightarrow s\gamma(g)$ :  
(use unitarity to factor out  $V_{tb}V_{ts}^*$ )

$$\mathcal{H}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_{i=1,2,7,8} C_i(\mu) O_i(\mu)$$

$$O_1 = (\bar{c}_{L\beta} \gamma^\mu b_{L\alpha})(\bar{s}_{L\alpha} \gamma_\mu b_{L\beta})$$

$$O_2 = (\bar{c}_{L\alpha} \gamma^\mu b_{L\alpha})(\bar{s}_{L\beta} \gamma_\mu b_{L\beta})$$

$$O_7 = \frac{em_b}{16\pi^2} \bar{s}_\alpha \sigma^{\mu\nu} (1 + \gamma_5) b_\alpha F_{\mu\nu}$$

$$O_8 = \frac{g_s m_b}{32\pi^2} \bar{s}_\alpha \sigma^{\mu\nu} (1 + \gamma_5) \lambda_{\alpha\beta}^a b_\beta G_{\mu\nu}^a$$

The lowest level:  $O_7$  only.

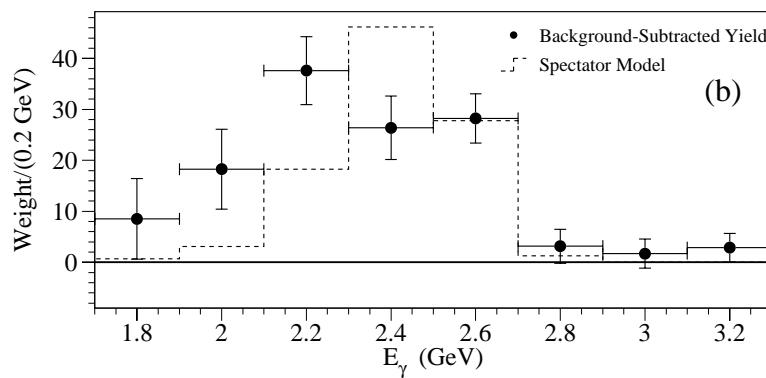
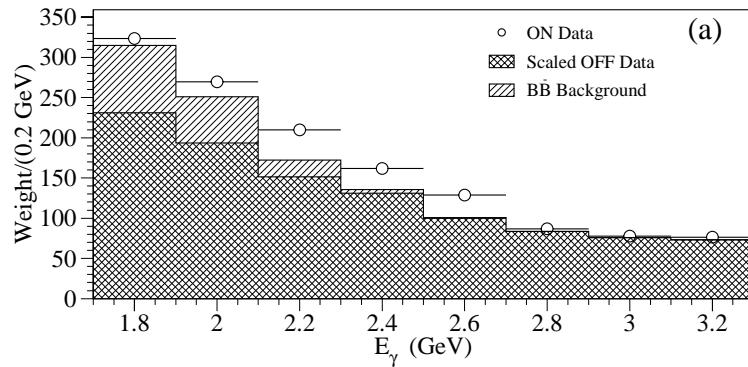
QCD correction enhances the rate by  $\sim 3$ .

- Complete NLO (order  $\alpha_s$ ) calculation done.
- Scale dependence  $\sim 20\%$  (LO)  $\rightarrow \sim 5\%$  (NLO).
- ‘Non perturbative effects’ ( $\lambda_{1,2}$ )  $\sim +3\%$ .

$b \rightarrow s\gamma$  CLEO ( $3 \text{ fb}^{-1}$ )

Look for an excess of single photon at the  $B$  decay end point ( $2.1 < E_\gamma < 2.7 \text{ GeV}$ )

Suppress continuum bkg by ‘pseudo  $B$  reconstruction’:  $\gamma + 1K^{0,+} + n^+\pi^\pm + n^0\pi^0$  ( $n^+ + n^0 \leq 4, n^0 \leq 1$ ) that is consistent with  $B$



Dotted line: shape predicted for  $m_b = 4.88 \text{ GeV}$  and Fermi motion  $p_B = 250 \text{ MeV}$  (model dep.)

## CLEO

$$Br(b \rightarrow s\gamma) = (3.15 \pm 0.35 \pm 0.32 \pm 0.26) \times 10^{-4}$$

stat    sys    model

## ALEPH

$$Br(b \rightarrow s\gamma) = (3.11 \pm 0.80 \pm 0.72) \times 10^{-4}$$

## Theory

$$Br(b \rightarrow s\gamma) = (3.28 \pm 0.33) \times 10^{-4}$$

Agreement is excellent.

'Pseudo  $B$  reconstruction'  $\rightarrow b$  flavor tag.

Mistag rate  $\sim 8\%$  (CLEO)

$$-0.09 < A_{CP} \equiv \frac{b - \bar{b}}{b + \bar{b}} < 0.42$$

$$\text{or } A_{CP} = (+0.16 \pm 0.14 \pm 0.05) \times (1 \pm 0.04)$$

mistag

Prospect: CLEO has  $3\times$  more data

$\rightarrow$  soon to produce a result.

**Exclusive**  $B \rightarrow X_{s,d}\gamma$   
 (CLEO)

- $B^+ \rightarrow K^{*+}\gamma, B^0 \rightarrow K^{*0}\gamma$

$$\begin{aligned} K^{*+} &\rightarrow K^+\pi^0, K^0\pi^+ \\ K^{*0} &\rightarrow K^+\pi^-, K^0\pi^0 \end{aligned}$$

$$\frac{B \rightarrow K^*\gamma}{B \rightarrow X_s\gamma}, \quad A_{CP}(\text{small in SM}).$$

- $B \rightarrow K_2^*(1430)\gamma$

The same final states as  $K^*$  above.

- $B \rightarrow (\rho/\omega)\gamma$

$$\frac{Br(B \rightarrow (\rho/\omega)\gamma)}{B(B \rightarrow K^*\gamma)} = \xi \left| \frac{V_{td}}{V_{ts}} \right|^2$$

$\xi$ : SU(3) breaking, phase space, long-distance

## <Full reconstruction on $\Upsilon 4S$ >

$$B \rightarrow f_1 \cdots f_n$$

Energy and absolute momentum of  $B$  known:

$$E_B = E_{\text{beam}} = 5.290 \text{ GeV}$$
$$|\vec{P}_B| = \sqrt{E_{\text{beam}}^2 - M_B^2} = 0.34 \text{ GeV/c}$$

→ require that candidates satisfy

$$E_{\text{tot}} = E_{\text{beam}}, \quad |\vec{P}_{\text{tot}}| = |\vec{P}_B|$$

where

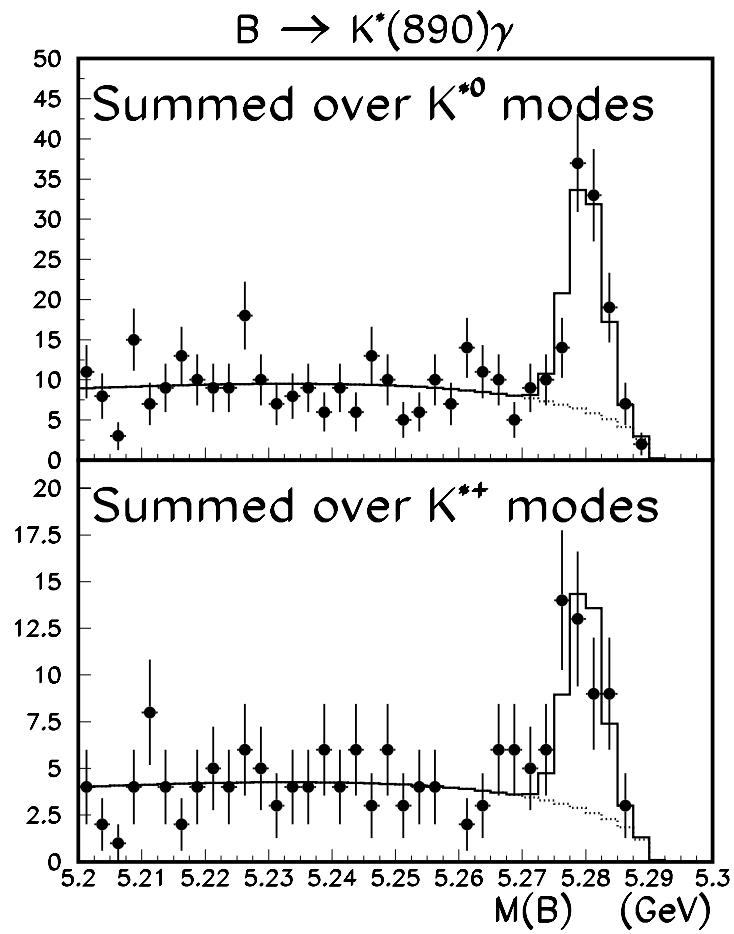
$$E_{\text{tot}} \equiv \sum_{i=1}^n E_i, \quad \vec{P}_{\text{tot}} \equiv \sum_{i=1}^n \vec{P}_i$$

Instead of  $E_{\text{tot}}$  and  $|\vec{P}_{\text{tot}}|$ , we historically use

$$\Delta E \equiv E_{\text{tot}} - E_B \quad (\text{energy difference})$$

$$M_{\text{bc}} \equiv \sqrt{E_{\text{beam}}^2 - \vec{P}_{\text{tot}}^2} \quad (\text{beam-constrained mass})$$

## $B \rightarrow K^*\gamma$ (CLEO)

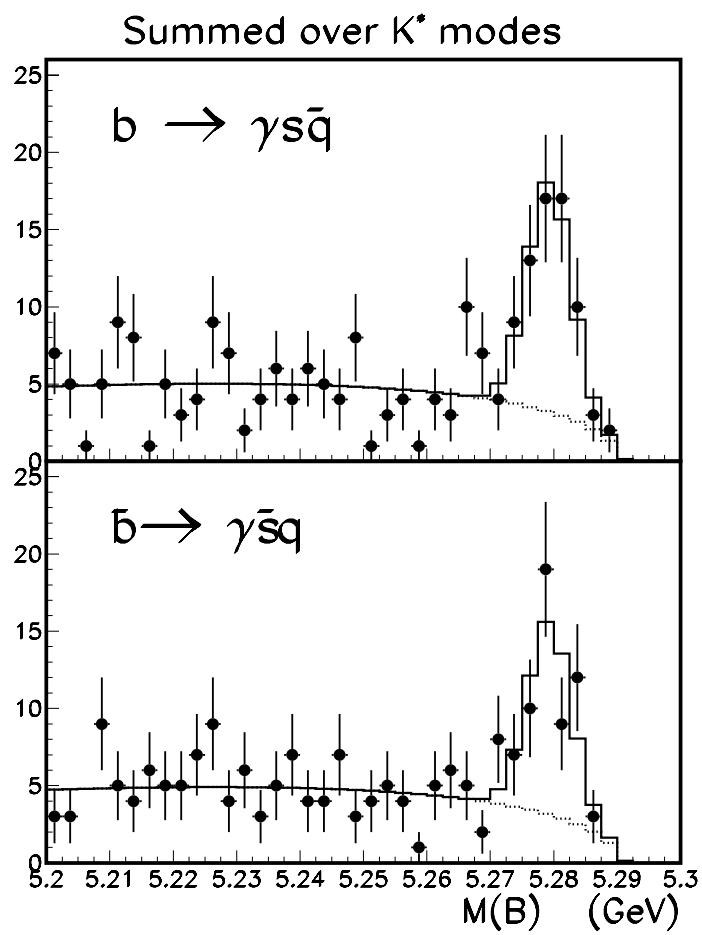


$$Br(B^0 \rightarrow K^{*0}\gamma) = (4.550_{-0.683}^{+0.715} \pm 0.343) \times 10^{-5}$$

$$Br(B^+ \rightarrow K^{*+}\gamma) = (3.764_{-0.831}^{+0.893} \pm 0.284) \times 10^{-5}$$

$B \rightarrow K^* \gamma$   $A_{CP}$  (**CLEO**)

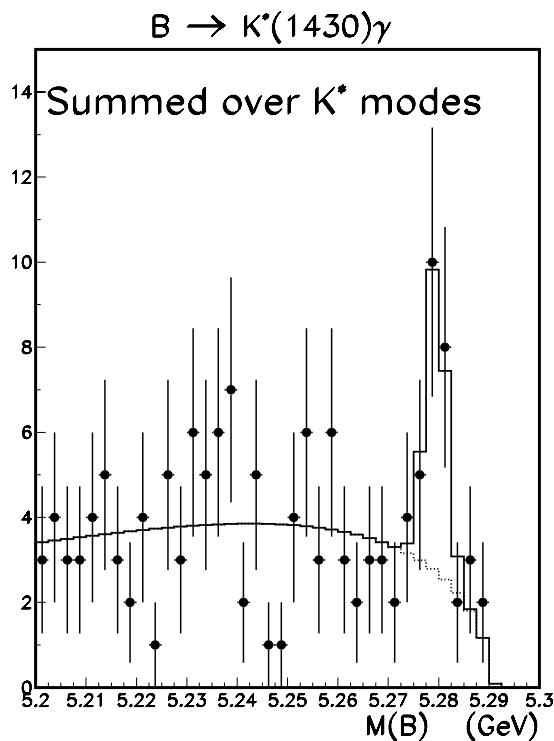
All modes other than  $K^0 \pi^0$  are flavor-tagged.



$$A_{CP} = +0.080 \pm 0.125 \quad (SM \sim 1\%)$$

$B \rightarrow K_2^*(1430)\gamma$  (**CLEO**)

$Br(K_2^*(1430)(J=2) \rightarrow K\pi) \sim 50\%$



$K^*(1410)(J=1)$  not excluded.

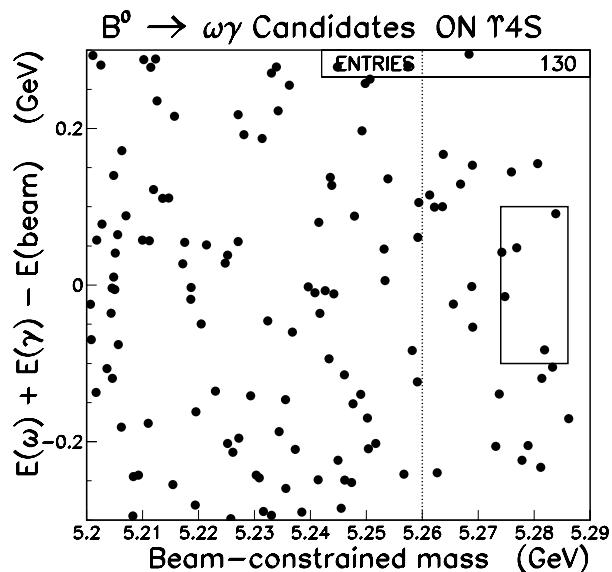
Assuming all  $K_2^*(1430)$ ,

$$Br(B \rightarrow K_2^*(1430)\gamma) = (1.7 \pm 0.5 \pm 0.1) \times 10^{-5}$$

$$R \equiv K_2^*(1430)\gamma / K^*(890)\gamma = 0.4 \pm 0.1$$

$$R = \begin{cases} 3.0 - 4.9 & (\text{Ali, Ohl, Mannel 93}) \\ 0.4 \pm 0.2 & (\text{Veseli, Olsson 96}) \end{cases}$$

## $B \rightarrow (\rho\omega)\gamma$ (**CLEO**)



Upper limits (90% c.l. 1996)

$$\begin{aligned} Br(B^0 \rightarrow \omega\gamma) &< 1 \times 10^{-5} \\ Br(B^0 \rightarrow \rho^0\gamma) &< 4 \times 10^{-5} \\ Br(B^- \rightarrow \rho^-\gamma) &< 1 \times 10^{-5} \end{aligned}$$

$$\rightarrow |V_{td}/V_{ts}|^2 < 0.45 \sim 0.56$$

Updated limit will come soon (3.5 times data).

# Unitarity Triangle

## Standard-Model quark-W Interaction

$$H_{\text{int}}(t) = \int d^3x [\mathcal{H}_{qW}(x) + \mathcal{H}_{qW}^\dagger(x)]$$

$$\mathcal{H}_{qW}(x) = \frac{g}{\sqrt{2}}(\bar{u}, \bar{c}, \bar{t})_L V_{CKM} \gamma_\mu \begin{pmatrix} d \\ s \\ b \end{pmatrix}_L W^\mu$$

$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \quad \begin{pmatrix} \text{CKM matrix} \\ \text{unitary} \end{pmatrix}$$

If one can adjust  $CP$  phases of quarks such that

$$(CP)\mathcal{H}_{qW}(x)(CP)^\dagger = \mathcal{H}_{qW}^\dagger(x') \quad (x' = (t, -\vec{x}))$$

$$\rightarrow (CP)H_{\text{int}}(t)(CP)^\dagger = H_{\text{int}}(t)$$

$$S = \sum_n \frac{(-i)^n}{n!} \int dt_1 \cdots dt_n T(H_{\text{int}}(t_1) \cdots H_{\text{int}}(t_n))$$

$$\rightarrow (CP)S(CP)^\dagger = S$$

- Then, the transition amplitude of  $|i\rangle \rightarrow |f\rangle$  becomes the same as that of  $CP|i\rangle \rightarrow CP|f\rangle$ :

$$\langle f|S|i\rangle = \langle f|(CP)^\dagger \underbrace{(CP)S(CP)^\dagger}_{S}(CP)|i\rangle = \langle f|(CP)^\dagger S(CP)|i\rangle$$

I.e., same amp. if particle↔antiparticle and helicity flip in the initial and final states.

- Or, if  $|i\rangle$  and  $|f\rangle$  are eigen states of  $CP$ :

$$CP|i\rangle = \eta_i|i\rangle, \quad CP|f\rangle = \eta_f|f\rangle,$$

Then,

$$\langle f|S|i\rangle = \eta_f^* \eta_i \langle f|S|i\rangle;$$

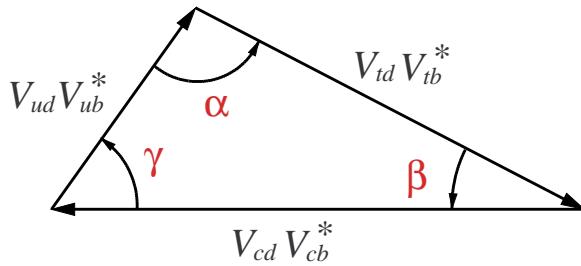
I.e., no transition unless  $\eta_f^* \eta_i = 1$  or  $\eta_f = \eta_i$ .  
CP eigenvalue is conserved.

In general, a  $3 \times 3$  unitary matrix has one non-trivial phase that cannot be rotated away by adjusting the  $CP$  phases of quarks. → CPV.

## A Main Question of the CPV Study in B: ‘Is $V_{CKM}$ unitary?’

e.g: orthogonality of  $d$ -column and  $b$ -column:

$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$



$$\begin{aligned}\alpha(\phi_2) &\equiv \arg \left( \frac{V_{td}V_{tb}^*}{-V_{ud}V_{ub}^*} \right), \beta(\phi_1) \equiv \arg \left( \frac{V_{cd}V_{cb}^*}{-V_{td}V_{tb}^*} \right) \\ \gamma(\phi_3) &\equiv \arg \left( \frac{V_{ud}V_{ub}^*}{-V_{cd}V_{cb}^*} \right)\end{aligned}$$

Note: the definitions of  $\alpha, \beta, \gamma$  are independent of quark phases.

For **any** 3 complex numbers  $a, b, c$ , by defintion,

$$\arg\left(\frac{c}{-a}\right) + \arg\left(\frac{b}{-c}\right) + \arg\left(\frac{a}{-b}\right) = \pi \pmod{2\pi}$$

regardless of  $a + b + c = 0$  or not.

It does not test the unitarity.

It tests angles measured are as defined or not.  
(e.g. due to penguin contamination etc.)

Need also to measure the lengths of the sides for unitarity test.

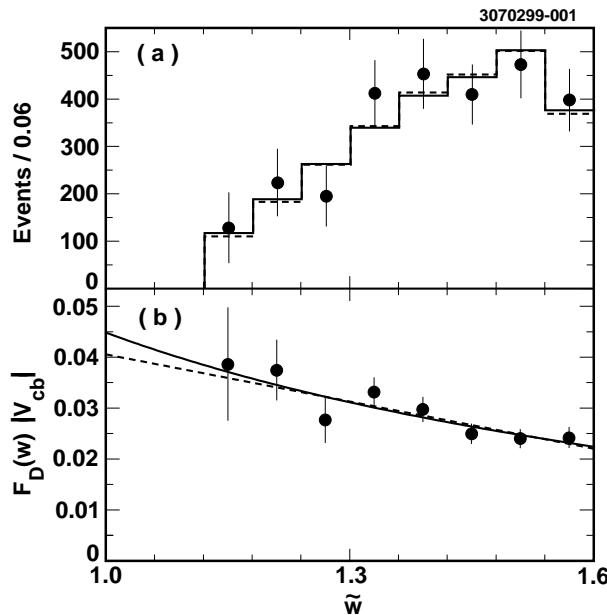
## Measurement of $|V_{cb}|$ (CLEO) (1999)

$$B \rightarrow D^{0,+} \ell \nu$$

$$\frac{d\Gamma}{dw} = \frac{G_F^2 |V_{cb}|^2 F_D(w)^2}{48\pi^3} (m_B + m_D)^2 m_D^3 (w^2 - 1)^{3/2}$$

$w$ :  $\gamma$ -factor of  $D$  in  $B$  frame.

$F_D(1) = 1$  in the HQ limit.



$$Br(\bar{B}^0 \rightarrow D^+ \ell \nu) = 2.09 \pm 0.13 \pm 0.18\% \\ Br(B^- \rightarrow D^0 \ell \nu) = 2.21 \pm 0.13 \pm 0.19\%$$

$$|V_{cb}| F_D(1) = 0.0416 \pm 0.0047 \pm 0.0037 \\ |V_{cb}| = 0.042 \pm 0.005 \pm 0.004 \pm 0.004$$

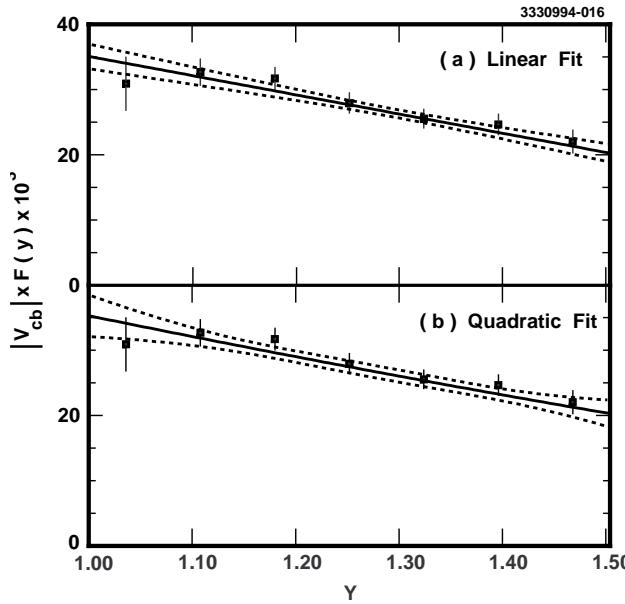
$$(F_D(1)=1.0 \pm 0.1)$$

$B \rightarrow D^{*0,+} \ell \nu$  (CLEO) (1995)

$$\frac{d\Gamma}{dw} = \frac{G_F^2 |V_{cb}|^2 F(w)^2}{48\pi^3} (m_B + m_{D^*})^2 m_D^{*3} \times \sqrt{w^2 - 1} \left[ 4w(w+1) \frac{1 - 2wr + r^2}{(1-w)^2} + (w+1)^2 \right]$$

$$r = m_{D^*}/m_B$$

F(1) = 1 in the HQ limit.



$$|V_{cb}|F(1) = 0.0351 \pm 0.0019 \pm 0.0018$$

$$|V_{cb}| = 0.0383 \pm 0.0021 \pm 0.0020 \pm 0.0011$$

New preliminary Br's:

$$Br(\bar{B}^0 \rightarrow D^{*+} \ell \nu) = 4.91 \pm 0.17 \pm 0.42\%$$

$$Br(B^- \rightarrow D^{*0} \ell \nu) = 5.13 \pm 0.54 \pm 0.64\%$$

New DELPHI result ON  $B \rightarrow D^* \ell \nu$  (1999)

$$Br(B^0 \rightarrow D^* \ell \nu) = 5.22 \pm 0.12 \pm 0.55\%$$

$$|V_{cb}|A(1) = 0.0380 \pm 0.0013 \pm 0.0016$$

$$|V_{cb}| = 0.0417 \pm 0.0015 \pm 0.0017 \pm 0.0014$$

$(A_1(1) \sim F(1) = 0.91 \pm 0.03)$

LEP combined result on  $|V_{cb}|$ :

Inclusive  $(40.8 \pm 0.4 \pm 2.0) \times 10^{-3}$

Exclusive  $(38.4 \pm 2.5 \pm 2.2) \times 10^{-3}$

average  $(40.2 \pm 1.9) \times 10^{-3}$

- The agreement between exclusive and inclusive indicates validity of global quark-hadron duality.
- $D^* \ell \nu$  has more yield near  $w = 1$  than  $D \ell \nu$   
→ less model dependence.

## Inclusive $b \rightarrow u\ell\nu$ measurements at LEP

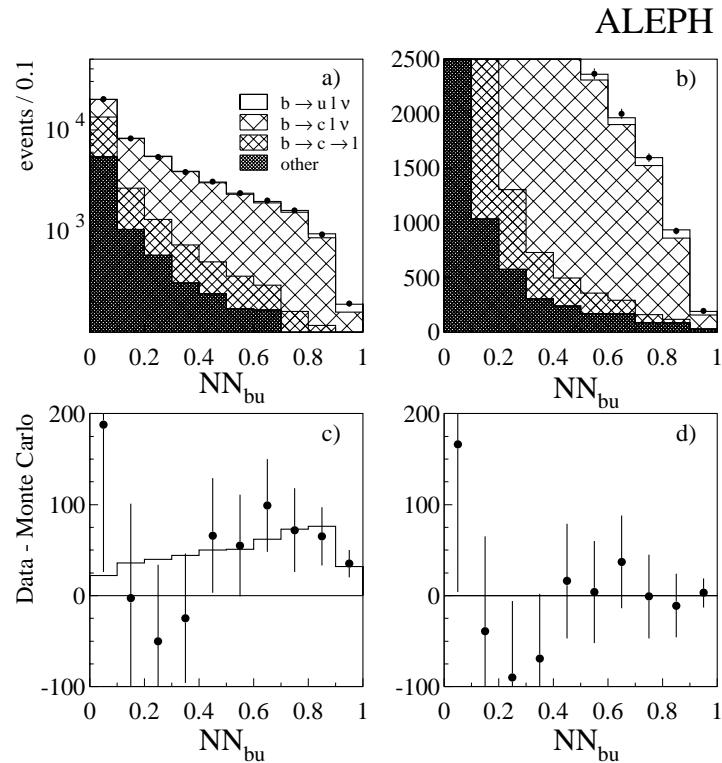
ALEPH, DELPHI, L3 combined

$$Br(b \rightarrow u\ell\nu) = (1.67 \pm 0.35 \pm 0.38 \pm 0.20) \times 10^{-3}$$

model

$$|V_{ub}| = \left( 4.05^{+0.39+0.43+0.23}_{-0.46-0.51-0.27} \pm 0.02 \pm 0.16 \right) \times 10^{-3}$$

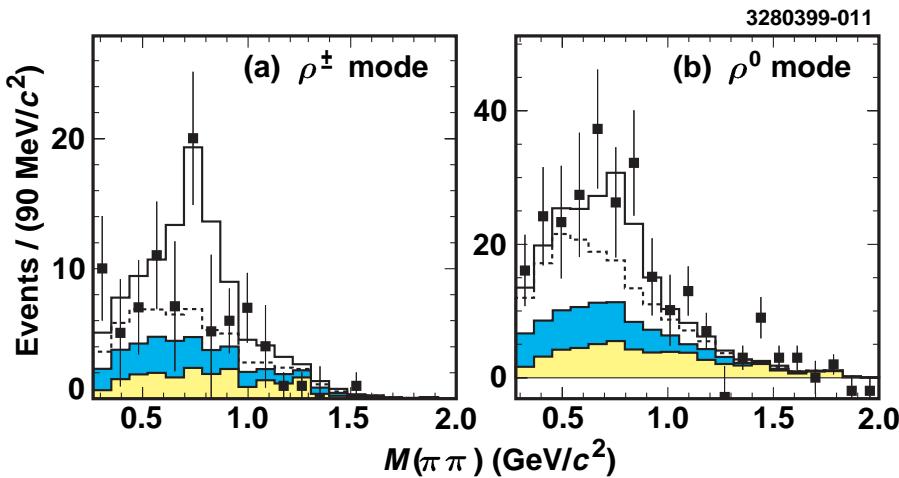
$\tau_B$               Theory



$NN_{bu}$ : neural net variable  
(even shape, charged multiplicity, invariant mass etc.)

## Exclusive $B \rightarrow \rho^- \ell \nu$ measurements by CLEO

Sensitivity is at high lepton energy  
( $E_\ell > 2.3$  GeV)



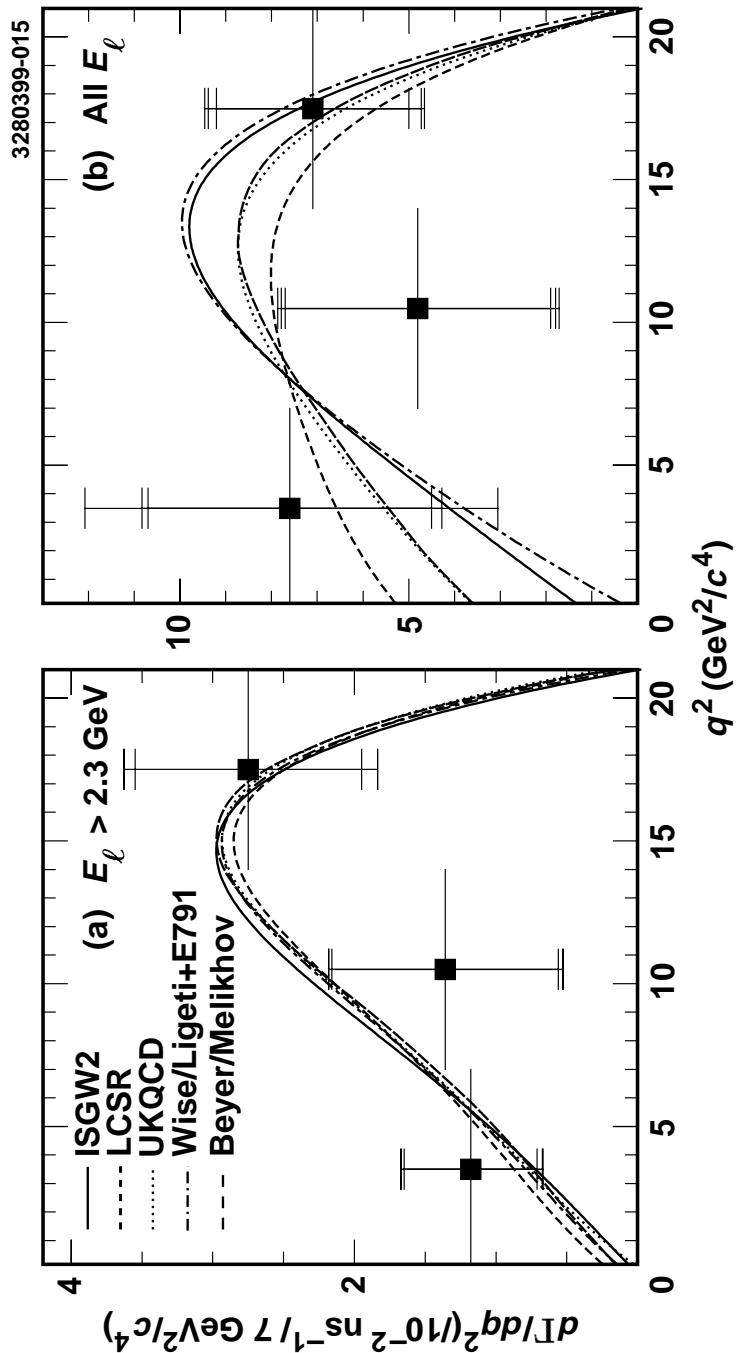
Combining with the old results

$$Br(B \rightarrow \rho^- \ell \nu) = (2.57 \pm 0.29^{+0.33}_{-0.46} \pm 0.41) \times 10^{-4}$$

$$|V_{ub}| = (3.25 \pm 0.14^{+0.21}_{-0.29} \pm 0.55(\text{th.})) \times 10^{-3}$$

$$(\rho \ell \nu + \pi \ell \nu)$$

CLEO  $B \rightarrow \rho \ell \nu$



## $B$ -mixing ( $|V_{td}|$ )

$$\chi \equiv \frac{B(t=0) \rightarrow \text{bar}B(\text{at decay})}{B(t=0) \rightarrow \text{all}} = \frac{1}{2} \frac{x^2}{1+x^2}$$

$$x \equiv \frac{\Delta m}{\Gamma}, \quad (\Gamma = (1.6 ps)^{-1})$$

$$\Delta m_d = \frac{g^4}{192\pi^2 m_W^2} |V_{tb}|^2 |V_{td}|^2 f_{B_d}^2 B_{B_d}^2 \eta_B f(m_t)$$

- Measure the mixing rate  $\chi \rightarrow |V_{td}|$ .  
Large theoretical uncertainties ( $f_{B_d}$  and  $B_{B_d}$ )

- Or, measure  $B_d$  mixing and  $B_s$  mixing:

$$\frac{\Delta m_s}{\Delta m_d} = \frac{m_{B_s}}{m_{B_d}} \left| \frac{V_{ts}}{V_{td}} \right|^2 \xi^2 \frac{\eta_{B_s}}{\eta_{B_d}}$$

Less theoretical uncertainty, but

$\Delta m_s$  not measured yet:

$$\Delta m_s > 14.3 \text{ ps}^{-1}$$

## A New CLEO Measurement of $B_d$ Mixing (Preliminary)

On  $\Upsilon 4S$ :

$$\chi = \frac{\text{Both decay as } B + \text{Both decay as } \bar{B}}{\text{All}}$$

- One side: lepton tag
- The other size: partial reconstruction of

$$B \rightarrow D^{*+} n\pi, D^{*+} \rightarrow D^0 \pi^+$$

$\vec{v}(D^*) \sim \vec{v}(\pi^+)$  allows the reconstruction of  $B \rightarrow D^{*+} n\pi$  without  $D^0$ .

	mixed	unmixed
$e$ -tag	$210 \pm 17$	$838 \pm 30$
$\mu$ -tag	$191 \pm 16$	$626 \pm 26$

Correct for mistags

$$\chi_d = 19.7 \pm 1.3\%$$